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Buckling of a Cylindrical Shell Loaded by a Pre-Tensioned Filament Winding

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Nomenclature

D	=	plate flexural stiffness, $Et^3/12(1-\mu^2)$
\boldsymbol{E}	=	Young's modulus
\boldsymbol{k}	=	elastic spring constant
k_y	=	buckling coefficient, $NL^2/D\pi^2$
$ {K}$	=	spring stiffness parameter, kL^4/π^4D
L	=	length of cylinder
n	=	number of waves in circumferential direction
N	=	applied circumferential stress resultant
N_x, N_y, N_{xy}	=	in-plane stress resultants
r	=	radius of cylindrical shell
t		thickness of cylindrical shell
		dianta coments in the m as and radial direct
u,v,w	=	displacements in the x-, y-, and radial direc-
u,v,w	=	displacements in the x-, y-, and radial directions, respectively
u,v,w x,y		tions, respectively
	-	tions, respectively axial and circumferential coordinates, respectively
x,y Z	=	tions, respectively axial and circumferential coordinates, respec-
x,y Z $\epsilon_x,\epsilon_y,\gamma_{xy}$	=======================================	tions, respectively axial and circumferential coordinates, respectively curvature parameter, $(L^2/rt)(1-\mu^2)^{1/2}$
x,y Z	=======================================	tions, respectively axial and circumferential coordinates, respectively curvature parameter, $(L^2/rt)(1-\mu^2)^{1/2}$ in-plane strains
x,y Z $\epsilon_x,\epsilon_y,\gamma_{xy}$ μ	=======================================	tions, respectively axial and circumferential coordinates, respectively curvature parameter, $(L^2/rt)(1-\mu^2)^{1/2}$ in-plane strains Poisson's ratio
x,y Z $\epsilon_x, \epsilon_y, \gamma_{xy}$ μ ∇^4	=======================================	tions, respectively axial and circumferential coordinates, respectively curvature parameter, $(L^2/rt)(1-\mu^2)^{1/2}$ in-plane strains Poisson's ratio $\partial^4/\partial x^4 + 2(\partial^4/\partial x^2\partial y^2) + \partial^4/\partial y^4$
x,y Z $\epsilon_x, \epsilon_y, \gamma_{xy}$ μ ∇^4		tions, respectively axial and circumferential coordinates, respectively curvature parameter, $(L^2/rt)(1-\mu^2)^{1/2}$ in-plane strains Poisson's ratio $\partial^4/\partial x^4 + 2(\partial^4/\partial x^2\partial y^2) + \partial^4/\partial y^4$ inverse operator defined by equation $\nabla^{-4}(\nabla^4 f) =$

When the subscripts x and y follow a comma, they indicate partial differentiation with respect to x and y, respectively.

Introduction

DESIGNS have appeared recently in launch-vehicle technology in which circular-cylindrical shells are wrapped with pre-tensioned filaments. An example of such a design is a

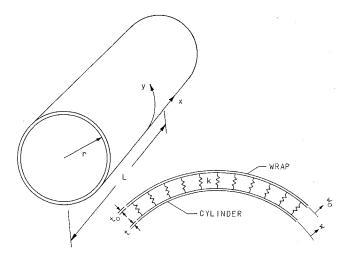


Fig. 1 Configuration and coordinate system.

cylinder that is wrapped circumferentially to help carry hoop stress from internal pressure. Another example is a cylindrical cryogenic tank upon which insulation is held in place by pre-tensioned filaments to protect against aerodynamic heating. One question that arises in the design of such configurations is: What is the largest filament tension that can be allowed before buckling of the shell would occur? It is believed that the theoretical analysis presented in this note is a contribution toward answering this question.

An idealized model of the configuration is studied in which material compressibility in the thickness direction is represented by many closely spaced, independent, radial springs located between the shell and the filaments. For the case of insulation secured to the shell by filament winding, the springs represent essentially the thickness compressibility of the insulation. For the case of filaments wrapped directly on the shell, the springs represent some effective thickness compressibility of the shell and filaments. A simple, closed-form solution is obtained for the idealized model. Results are plotted in dimensionless form so that the buckling load (the critical tension in the filaments or the critical compression in the shell) can be obtained directly for wide ranges of the geometrical and stiffness parameters of the model.

Analysis

The buckling of an elastic isotropic, homogeneous, circular-cylindrical shell loaded through many closely spaced, radial springs by pre-tensioned circumferentially wrapped filaments is analyzed (see Fig. 1). The filament winding is considered to have no bending stiffness. The difference between the radius of the shell and of the filaments is assumed to be negligible so that the hoop tension per unit length in the filaments N causes an equal magnitude of hoop compression -N in the shell. The filaments are permitted to slide freely relative to the cylinder and relative to each other; i.e., no bonding agent is applied to the filaments or the cylinder. Thus the hoop tension load per unit length in the filament winding is simply the number of filaments per unit length times the load in each filament.

For changes in stress resultants and displacements, which occur during buckling, the equilibrium equations are, for the shell,

$$N_{x,x} + N_{xy,y} = 0 (1)$$

$$N_{y,y} + N_{xy,x} = 0 (2)$$

$$D\nabla^4 w + (N_y/r) + Nw_{,yy} + k(w - w_0) = 0$$
 (3)

and, for the filament winding,

$$N_{y0,y} = 0 \tag{4}$$

Received September 29, 1964.

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$$(N_{y0}/r) - Nw_{0,yy} - k(w - w_0) = 0 (5)$$

The stress-resultant strain relations are, for the shell,

$$N_{x} = \frac{Et}{1 - \mu^{2}} (\epsilon_{x} + \mu \epsilon_{y})$$

$$N_{y} = \frac{Et}{1 - \mu^{2}} (\epsilon_{y} + \mu \epsilon_{x})$$

$$N_{xy} = \frac{Et}{2(1 + \mu)} \gamma_{xy}$$
(6)

and for the filament winding

$$N_{y0} = E_0 t_0 \epsilon_{y0} \tag{7}$$

The strain-displacement relations are, for the shell,

$$\begin{cases}
\epsilon_x = u_{,x} \\
\epsilon_y = v_{,y} + (w/r)
\end{cases}$$

$$\begin{cases}
\epsilon_y = u_{,y} + (w/r) \\
\epsilon_y = u_{,y} + v_{,x}
\end{cases}$$
(8)

and, for the filament winding

$$\epsilon_{y0} = v_{0,y} + (w_0/r)$$
 (9)

In a procedure similar to that used in Ref. 1, Eqs. (1-3, 6, and 8) can be combined into one equation for the lateral displacement of the shell w as

$$D\nabla^{4}w + \frac{Et}{r^{2}}\nabla^{-4}w_{,xxxx} + Nw_{,yy} + k(w - w_{0}) = 0 \quad (10)$$

Upon partial differentiation of Eq. (5) with respect to y and utilization of Eq. (4), the following equation is obtained:

$$Nw_0, y_{yy} + k(w - w_0)_y = 0 (11)$$

Equations (10) and (11) can be solved exactly for the classical simple support boundary conditions ($w = w_0 = 0$, w, $_{xx} = 0$; at 0, L) if the lateral displacements are chosen as $w = A \sin(\pi x/L) \sin(ny/r)$ and $w_0 = B \sin(\pi x/L) \sin(ny/r)$. Substitution of these expressions for the lateral displacements (w and w_0) into Eqs. (10) and (11) yields the following two homogeneous equations for A and B:

$$AD \left[\left(\frac{\pi}{L} \right)^2 + \left(\frac{n}{r} \right)^2 \right]^2 + \frac{A (Et/r^2) (\pi/L)^4}{\left[(\pi/L)^2 + (n/r)^2 \right]^2} - NA \left(\frac{n}{r} \right)^2 + k(A - B) = 0 \quad (12)$$

$$-NB\left(\frac{n}{r}\right)^2 + k(A - B) = 0 \tag{13}$$

By setting the determinant of the coefficients of A and B equal to zero, the following equation for the critical value of the stress resultant N can be obtained after some manipulation:

$$\frac{k_y^2\beta^2}{k_y\beta^2 + K} = \frac{(1+\beta^2)^2}{\beta^2} + \frac{12Z^2}{\pi^4} \frac{1}{(1+\beta^2)^2\beta^2}$$
(14)

where

$$Z = \frac{L^2}{rt} (1 - \mu^2)^{1/2} \qquad K = \frac{kL^4}{D\pi^4}$$
$$k_y = \frac{NL^2}{\pi^2 D} \qquad \beta = \frac{nL}{\pi r}$$

Equation (14) can be solved explicitly for k_y as

$$k_y = \frac{k_{yp} + [k_{yp}^2 + (4k_{yp}K/\beta^2)]^{1/2}}{2}$$
 (15)

where

$$k_{ur} = \frac{(1+\beta^2)^2}{\beta^2} + \frac{12Z^2}{\pi^4} \frac{1}{\beta^2(1+\beta^2)^2}$$

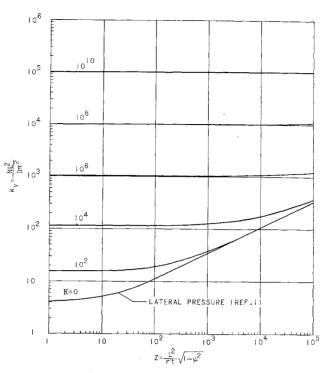


Fig. 2 Buckling coefficient for cylinder wrapped with pre-tensioned filaments.

is the equation from which the buckling characteristics of a cylinder loaded by lateral pressure may be determined.¹

Results

The critical stress resultant is determined by minimizing k_y in Eq. (15) with respect to β . The results for several values of the spring stiffness parameter K are given in terms of dimensionless parameters in Fig. 2. For the limiting case in which the spring stiffness parameter K approaches zero, the results are the same as those given in Ref. 1 for a cylinder loaded by lateral pressure. As the spring stiffness parameter approaches infinity, the buckling load becomes infinite. Note that no radial bond between the cylinder and springs is necessary to obtain this increase in the buckling load since the springs just prior to buckling are compressed.

For moderate Z and $K > 10^4$ it can be seen from Fig. 2 that a very good approximate formula for the buckling coefficient k_y can be written as

$$k_u = K^{1/2} (16)$$

or

$$N_{\rm er} = (Dk)^{1/2} \tag{17}$$

The quantity $N_{\rm cr}$ is the critical tension load per unit length in the wrap.

Discussion

The results of this analysis show that circular-cylindrical shells loaded by pre-tensioned filament windings can have much higher classical buckling strengths than would be expected if the loading were replaced by an equivalent external lateral pressure. The analysis demonstrates that, for initially perfect shells, material compressibility in the thickness direction plays a major role in the mechanism through which these structures buckle. For configurations that are very stiff in the thickness direction, other effects such as initial imperfections might be important.

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